

# Melting above the anhydrous solidus controls the location of volcanic arcs

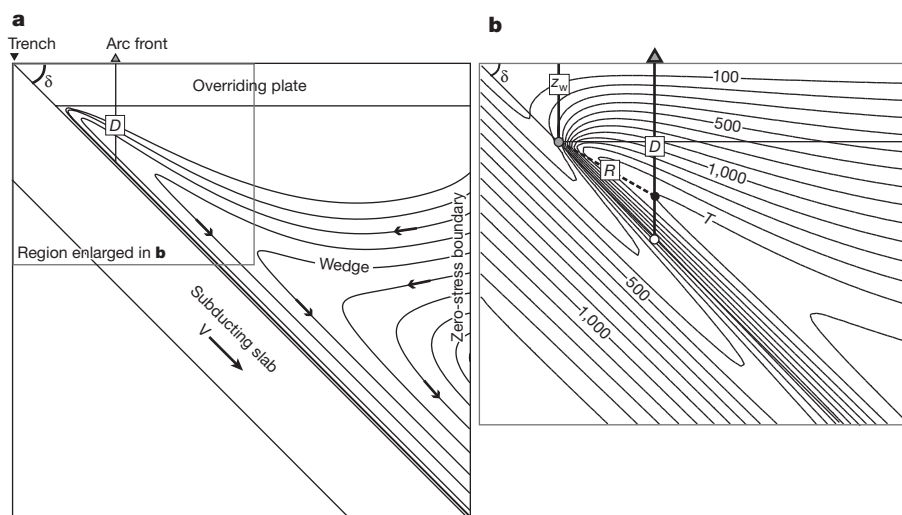
Philip C. England<sup>1</sup> & Richard F. Katz<sup>1</sup>

Segregation of magma from the mantle in subduction zones is one of the principal mechanisms for chemical differentiation of the Earth. Fundamental aspects of this system, in particular the processes by which melt forms and travels to the Earth's surface, remain obscure. Systematics in the location of volcanic arcs, the surface expression of this melting, are widely considered to be a clue to processes taking place at depth, but many mutually incompatible interpretations of this clue exist (for example, see refs 1–6). We discriminate between those interpretations by the use of a simple scaling argument derived from a realistic mathematical model of heat transfer in subduction zones. The locations of the arcs cannot be explained by the release of fluids in reactions taking place near the top of the slab. Instead, the sharpness of the volcanic fronts, together with the systematics of their locations, requires that arcs must be located above the place where the boundary defined by the anhydrous solidus makes its closest approach to the trench. We show that heat carried by magma rising from this region is sufficient to modify the thermal structure of the wedge and determine the pathway through which both wet and dry melts reach the surface.

Volcanic arcs are characterized by sharp fronts whose locations may be described, with misfits of no more than a few kilometres, by small circles on the Earth's surface (Fig. 1 and refs 7 and 8); furthermore, the depth of the top of the slab beneath these fronts falls in a narrow range

( $\sim 120 \pm 40$  km; refs 1, 3 and 5). The sharpness of the volcanic fronts implies that a key process in the generation or transport of magma is similarly focused beneath the arcs, but there is no consensus as to what that process may be. A wide range of metamorphic and melting reactions, either in the slab or in the mantle wedge, have been proposed as candidate processes. Some authors suggest that the arcs lie above places where the degree of melting in the mantle wedge becomes high enough for the melt to segregate from the solid<sup>2,4</sup>. Others suggest that the locations of arcs are determined by the release of fluid near the top of the slab in reactions that are either strongly pressure-dependent<sup>1,3,5</sup> or strongly temperature-dependent<sup>6,9</sup>. None of these suggestions has, however, produced a successful quantitative prediction of the location of volcanic arcs. Here we take a different approach: starting with the observed correlation between the descent speed of the slab and its depth beneath the volcanic arcs<sup>8</sup>, we use a simple mathematical model to fit the data and reveal the petrological processes responsible for the locations of the arcs.

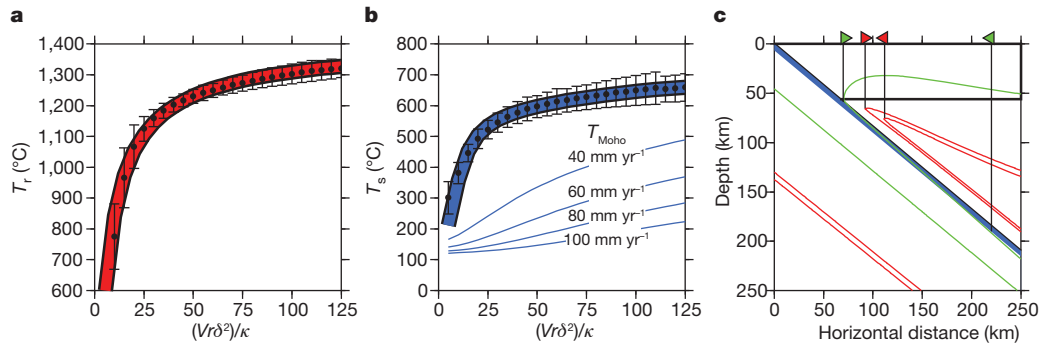
Although calculations of the full temperature field in subduction zones require numerical models<sup>10</sup>, their results can be encapsulated in simple scaling relations that show<sup>11</sup> that temperatures within the mantle wedge and at the top of the slab depend upon a single parameter,  $Vr\delta^2/\kappa$ . Here  $V$  is the convergence rate across the plate boundary,  $\delta$  is the dip of the slab,  $r$  is the radial distance from the wedge corner (Fig. 1), and  $\kappa$  is



**Figure 1** | Idealized cross-sections of a subduction zone, drawn perpendicular to the trench and the island arc. **a**, Two plates converge at a speed  $V$ , with the slab of oceanic lithosphere being subducted at an angle  $\delta$  beneath the overriding plate. The arc front is a zone a few kilometres wide, across which volcanic activity begins as one moves away from the trench; it lies at a distance  $D$  above the top of the slab. The relative motion between slab and overriding plate generates a creeping flow in the wedge of mantle between them, which follows stream lines, shown as curved lines. The corner of the

wedge is at a depth  $z_w$ . **b**, Enlargement of temperature structure of the rectangle in **a**: Isotherms are shown at intervals of  $100^\circ\text{C}$ . A schematic isotherm, labelled  $T$ , has its closest approach to the wedge corner (its 'nose') immediately beneath the volcanic front, at a distance  $R$  from the corner (black circle). An open circle marks the top of the slab directly below the corner of this isotherm, at a depth  $D$  below the surface. The distance  $R$  cannot be determined by observation, but is similar to  $R_D = (D - z_w)/\sin\delta$ , the radial distance of the top of the slab (open circle) from the wedge corner.

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**Figure 2 | Scaling relations for temperatures in the core of the mantle wedge, and at the top of the slab.** **a**, Maximum temperature in the wedge,  $T_r$ , as a function of dimensionless distance from the wedge corner,  $Vr\delta^2/\kappa$ , where  $V$  is convergence speed,  $\delta$  is dip of the slab, and  $\kappa$  is thermal diffusivity. Dots with error bars indicate the averages and standard deviations of  $T_r$ , determined from calculations in which  $V$  varies from  $10 \text{ mm yr}^{-1}$  to  $100 \text{ mm yr}^{-1}$ , in steps of  $10 \text{ mm yr}^{-1}$ , while  $\delta$  varies from  $20^\circ$  to  $70^\circ$  in steps of  $10^\circ$ . The red line corresponds to the theoretical expression for  $T_r$  (equation (1)), with  $T_0 = 1,420^\circ\text{C}$ ,  $B = 3.3$  and  $\beta = -0.8$ . **b**, As for **a**, but for temperatures within the slab. The thick blue line corresponds to the theoretical expression for  $T_s$ , the temperature on the top of the slab (equation (2)), with  $C = 1.4$  and  $\gamma = -0.06$ . Thin blue lines indicate the temperatures at the base of the crust (7 km below

the top of the slab) for convergence rates  $V$  of  $40\text{--}100 \text{ mm yr}^{-1}$  and a slab dip of  $40^\circ$ . **c**, The temperature structure near the wedge corner for a calculation with a convergence speed  $V$  of  $80 \text{ mm yr}^{-1}$  and a slab dip of  $40^\circ$ . The location of the  $500^\circ\text{C}$  isotherm is shown by green lines and the  $1,225^\circ\text{C}$  and  $1,275^\circ\text{C}$  isotherms are shown by red lines. Green arrowheads show the horizontal extent over which some part of the oceanic crust is at a temperature of  $500^\circ\text{C}$ . Red arrowheads show the range over which the maximum temperature in the mantle wedge lies between  $1,225^\circ\text{C}$  and  $1,275^\circ\text{C}$  (a typical range of temperature represented by the error bars in Fig. 3b). Uncertainties in the temperatures arising from idealizations in the model are discussed in section B2 of the Supplementary Information.

the thermal diffusivity of the mantle. In this scaling, the maximum temperature in the mantle wedge  $T_r$  is given by:

$$T_r \approx T_0 \exp \left[ -B \left( \frac{Vr\delta^2}{\kappa} \right)^\beta \right] \quad (1)$$

while the temperature  $T_s$  at the top of the slab is:

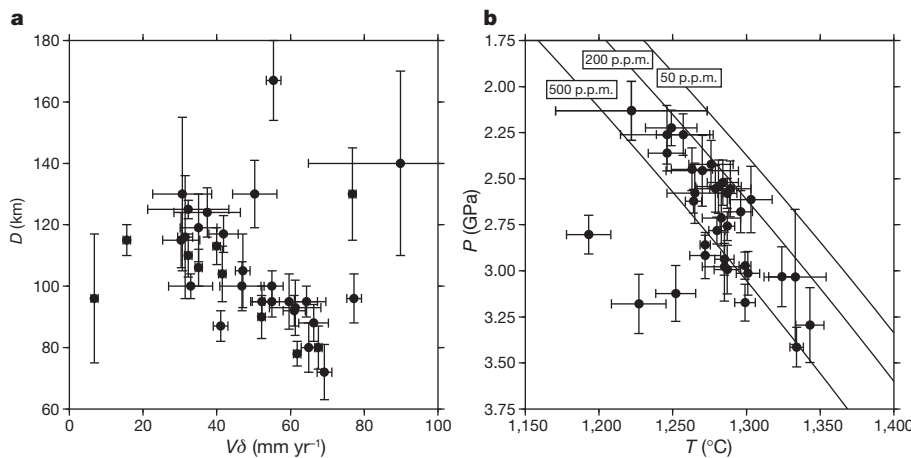
$$T_s \approx \frac{T_r}{1 + C(Vr\delta^2/\kappa)^\gamma} \quad (2)$$

where  $T_0$  is a scale temperature, and  $B$ ,  $C$ ,  $\beta$  and  $\gamma$  are constants, the values of which depend on the details of the flow near the top of the slab. This scaling was initially derived for a model of subduction zones that treated the mantle as a constant-viscosity fluid<sup>11</sup>; we show here that these relations also hold for the case in which the viscosity has a dependence on temperature that is appropriate for the upper mantle (Fig. 2 and Supplementary Information).

Precise earthquake hypocentral locations<sup>12</sup> reveal that the depth  $D$  to the top of the slab beneath the fronts of volcanic arcs is constant, to within a few kilometres, along individual segments of arc, but varies from 80 km to 160 km between different segments<sup>8,13</sup> (Fig. 3). This

variation rules out the hypothesis that the arcs are located above the place where the top of the slab reaches a critical pressure corresponding to a single dehydration reaction<sup>1,3,5</sup>. We may also rule out the hypothesis that the release of fluids by temperature-dependent reactions near the top of the slab determines the location of arcs (see, for example, ref. 6). The top of the slab lies within a thermal boundary layer a few tens of kilometres thick, across which there is a temperature difference of  $\sim 1,000^\circ\text{C}$ . Because isotherms within this boundary layer are almost parallel to the slab, any given temperature will be found over a large range of pressure (Fig. 2b, c). Therefore, temperature-dependent processes taking place near the top of the slab cannot be sharply localized, but must occur over a broad range of down-dip distances<sup>14</sup>.

In contrast, the steep lateral thermal gradients in the core of the mantle wedge provide a setting in which localization of temperature-dependent processes is likely<sup>2,4</sup>. The maximum temperature in the mantle wedge depends on the dimensionless distance from the wedge corner  $Vr\delta^2/\kappa$  (equation (1) and Fig. 2a). Accordingly, we should expect that if a temperature-dependent process is localized beneath the arc, the relevant isotherm will reach its closest approach to the wedge corner (which we refer to below as its ‘nose’) at a distance  $R$  that is inversely proportional to  $V\delta^2$ . Although  $R$  cannot be measured



**Figure 3 | Systematic variation in depth to the slab beneath volcanic arcs, and its relation to pressure-temperature conditions beneath the arcs.** **a**, Depth to the top of the slab beneath volcanic arcs (see Supplementary Information), plotted against  $V\delta$  (equation (3)). **b**, Conditions beneath the volcanic fronts estimated from calculations with descent speed,  $V$ , and dip,  $\delta$ , corresponding to the arcs investigated in this paper. Dots show the maximum temperature in the mantle wedge beneath the front, and the pressure at which that temperature is reached. For each calculation, the arc front is taken to lie immediately above the place where the top of the slab reaches the depth  $D$ ; error bars represent the range in maximum temperature (and the pressure at which it is reached) associated with a  $\pm 5 \text{ km}$  uncertainty in the horizontal location of the arc front. Lines labelled 50 p.p.m., 200 p.p.m. and 500 p.p.m. correspond to the solidi for peridotite containing these fractions of water<sup>15</sup>.

directly, it is very similar to the radial distance  $R_D = (D - z_w)/\sin(\delta)$  of the top of the slab from the wedge corner at that location (Fig. 1), thus:

$$\frac{VR_D\delta^2}{\kappa} \approx \frac{V\delta^2(D - z_w)}{\kappa \sin(\delta)} \approx \frac{V\delta(D - z_w)}{\kappa} \quad (3)$$

Here,  $z_w$  is the depth to the wedge corner, and we have made the small-angle approximation  $\sin(\delta) \approx \delta$ . Hence, for  $z_w$  constant or varying little in comparison with  $D$ , we should expect  $D$  to vary inversely with  $V\delta$ .

We can make a reliable determination of  $D$  for 35 arcs (Supplementary Information); the Spearman rank-order correlation coefficient between  $D$  and  $V\delta$  for these arcs (Fig. 3b) is  $-0.34$ , which is significant at the 95% level of confidence. The significance of the correlation is sensitive to the presence of five arcs with large uncertainties in  $D$  (identified in the discussion of Supplementary Fig. 1); the confidence level for the correlation exhibited when these five arcs are excluded is above 99%. Thirty arc segments, with a total length of 22,000 km, exhibit the relationship between  $D$  and  $V\delta$  that is predicted by equation (3), whereas five arc segments with a total length of  $\sim 4,000$  km do not lie on the trend. (See Supplementary Material for further discussion.) We therefore conclude that the global systematics<sup>8,13</sup> strongly support the hypothesis that the locations of volcanic arcs are determined by a temperature-dependent process taking place in the wedge<sup>2,4</sup>.

To estimate the conditions under which this process takes place, we carried out calculations of steady-state wedge thermal structure for combinations of slab dip and convergence rate corresponding to the 35 individual arcs in Fig. 3a. The maximum temperature in the column of mantle beneath the location of the arc, and the pressure at which that temperature is reached, form an array between about 1,250 °C and 1,325 °C, and 2–3.5 GPa (Fig. 3b) that corresponds closely to the range of pressure–temperature conditions for the melting of peridotite in the presence of between 200 p.p.m. H<sub>2</sub>O and 500 p.p.m. H<sub>2</sub>O (ref. 15). These water contents are about ten times lower than the concentrations estimated for wet melting beneath arcs, but are representative of the water contents of the mantle beneath backarc basins<sup>16</sup>. In what follows, we refer to such H<sub>2</sub>O concentrations as “anhydrous” to distinguish them from the hydrous melts that are much more abundant in the wedge.

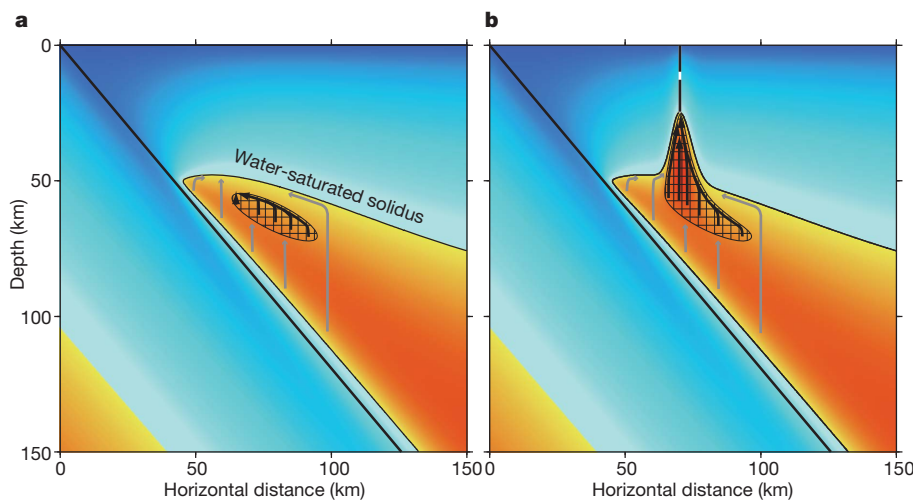
It is generally accepted that hydrous melting pervades the mantle wedge, but it has been suggested that the arc front is located somewhere

above the region in which the temperature in the wedge exceeds the anhydrous solidus, because the degree of hydrous melting increases rapidly there<sup>2,4</sup>. Melt may also be produced in this temperature range as a result of anhydrous decompression melting of upwelling mantle; indeed, the existence of distinct regimes of melting beneath the arcs has been inferred from the eruption of low-water-content, tholeiitic melt from arc volcanoes that also erupt typical hydrous melts<sup>17–21</sup>. Geodynamic models show, however, that either type of melting occupies a broad region of the wedge core, extending from beneath the arc front towards the back-arc<sup>22,23</sup>. Any explanation for the localization of the arc fronts that invokes melts produced over this broad volume must therefore include a mechanism for focusing the melts to the line beneath the arcs.

The interaction of rising magma with the thermal boundary layer at the top of the wedge provides this mechanism. As the rising magma hits this boundary, it begins to crystallize, forming an impermeable barrier above a sloping, high-porosity channel<sup>23,24</sup>. Such interaction requires sufficient crystallization to seal the pore space against vertical magmatic flow<sup>25</sup>; this condition is met at temperatures above the anhydrous solidus, where isobaric productivity is high and small decreases in temperature result in significant crystallization, but not below it, where isobaric productivity is low<sup>15,26</sup>. We therefore propose, as sketched in Fig. 4a, that the high-porosity channel terminates at the ‘nose’ of the anhydrous solidus. This ‘nose’ is controlled by temperature, so its location will depend upon  $V\delta$ , consistent with the observations (Fig. 3).

If the advective flux of heat carried by the magma is of the same order as the vertical conductive heat flux, the melt arriving at the nose of the solidus will erode the cold boundary layer at the top of the wedge and bow upward the solidi for both wet and dry melting. The rate of heat transfer by the melt is  $\sim ML$ , where  $M$  is the mass flow of magma and  $L$  is its latent heat per unit mass. The conductive heat flux is  $\sim k\Delta T/h$ , where  $k$  is thermal conductivity,  $\Delta T$  is the temperature difference between the solidus and the surface temperature, and  $h$  is the depth to the ‘nose’ of the solidus beneath the arc. This argument, expanded in the Supplementary Information, shows that the ratio of these two quantities controls the efficiency of thermal erosion:

$$\mathcal{M} = \frac{MLh}{k\Delta T} \quad (4)$$



**Figure 4 | A sketch of the process that determines the position of volcanoes.** The top of the slab is indicated by the diagonal line starting at the origin. The black line separating the blue and yellow-to-red colours within the wedge represents the water-saturated solidus of the mantle, the cross-hatched region is above the nominally anhydrous solidus of ambient mantle. **a**, The distribution of temperature and melting without heat transport by migrating melt. Grey and black arrows show melt formed above the water-saturated and anhydrous solidi, respectively, rising, then travelling through high-porosity channels to the ‘noses’ of the solidi. **b**, A schematic depiction of how melt transport, indicated

by arrows, would modify the distribution in **a**. Magma rising from the ‘nose’ of the dry solidus heats the region immediately above, deflecting upwards the dry and wet solidi so that they both reach their shallowest depth in the same horizontal location. Each type of melt travels laterally and upward along its respective solidus towards this location. Melts eventually penetrate the lithosphere by hydrofracture and dyking. In the case of low permeability in the wedge, advection of melt by the moving mantle can cause the trajectories to deviate horizontally<sup>23</sup>; this process could affect the details of the sketch.

For reasonable estimates of the rate of mass flux into arc volcanoes<sup>22,27–30</sup>  $\mathcal{M} \approx 0.8$ , which is sufficient to perturb upward the location of the anhydrous solidus by >40% of its depth (Supplementary Fig. 4). In consequence, all melt migrates towards the place where the solidus is bowed upward—including melt generated closer to the trench (Fig. 4b).

In summary, the observed systematics in the location of the volcanic arcs shows that the mechanisms whereby the melt reaches the Earth's surface are controlled by the dynamics of the wedge, not by any locally concentrated release of fluid in reactions taking place near the top of the slab. In our model, magma formed at temperatures above the anhydrous solidus in the wedge is focused trench-ward, to the 'nose' of the region bounded by that solidus. Above the 'nose', thermal erosion by rising magma establishes the pathway that all melts, hydrous or anhydrous, take to reach the arc volcanoes. Future analysis, including calculations that incorporate the physics of melt migration<sup>31</sup>, will help to clarify this hypothesis.

## METHODS SUMMARY

**Calculations of thermal structure of subduction zones.** Our calculation of the thermal structure of subduction zones employs a finite-volume discretization of the incompressible, variable-viscosity Stokes equation and steady-state, advection-diffusion energy equation on a uniformly spaced, staggered Cartesian mesh. It uses a Newton–Krylov solver provided by the Portable, Extensible Toolkit for Scientific Computation (<http://www.mcs.anl.gov/petsc>). The code has been benchmarked against other numerical solutions for this type of problem. The viscosity  $\eta$  is given by an Arrhenius law for diffusion creep of olivine:

$$\eta = A \exp\left(\frac{E + PV^*}{RT}\right) \quad (5)$$

with activation energy  $E$  and activation volume  $V^*$  of  $375 \text{ kJ mol}^{-1}$  and  $5 \times 10^{-6} \text{ m}^3 \text{ mol}^{-1}$ , respectively, and pre-exponential  $A = 1.8 \times 10^7 \text{ Pa s}$ .

The domain is 600 km by 600 km, oriented as shown in Fig. 1, with a mesh spacing of 1 km. We have checked convergence of our solutions by comparing them, for combinations of parameters that span the range of interest, against solutions with mesh spacing of 0.5 km. We find differences between the two meshes at the level of less than  $10^\circ \text{C}$  between temperatures calculated on the top of the slab or near the wedge corner.

**Measurement of depth to top of slab beneath arcs.** We apply the methods of ref. 8 to the 45 arc segments studied by ref. 13. For each arc segment we formed cross-sections of the seismicity and estimated  $D$  as described by ref. 8. We used the tables given in the supplementary information of ref. 13 to check whether the segment boundaries of the two studies agreed; except where noted in our Supplementary Table, the disagreements were minor and we used the segment boundaries of ref. 8.

**Full Methods** and any associated references are available in the online version of the paper at [www.nature.com/nature](http://www.nature.com/nature).

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**Supplementary Information** is linked to the online version of the paper at [www.nature.com/nature](http://www.nature.com/nature).

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**Author Contributions** R.F.K. wrote the code for the numerical experiments; P.C.E. carried out the re-analysis of depth-to-slab (Fig. 3 and Supplementary Information). Both authors participated equally in developing the ideas presented in this paper and in writing it.

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## METHODS

**Calculations of thermal structure of subduction zones.** Our calculation of the thermal structure of subduction zones employs a finite-volume discretization of the incompressible, variable-viscosity Stokes equation and steady-state, advection-diffusion energy equation on a uniformly spaced, staggered Cartesian mesh<sup>32</sup>. It uses a Newton–Krylov solver provided by the Portable, Extensible Toolkit for Scientific Computation<sup>33,34</sup>. The code has been benchmarked against other numerical solutions for this type of problem<sup>10</sup>. The viscosity  $\eta$  is given by an Arrhenius law for diffusion creep of olivine:

$$\eta = A \exp\left(\frac{E + PV^*}{RT}\right) \quad (6)$$

with activation energy  $E$  and activation volume  $V^*$  of  $375 \text{ kJ mol}^{-1}$  and  $5 \times 10^{-6} \text{ m}^3 \text{ mol}^{-1}$ , respectively, and pre-exponential  $A = 1.8 \times 10^7 \text{ Pa s}$  (refs 35 and 36).

The domain is composed of three subdomains, the slab, the wedge, and the overlying lid. Velocity in the slab subdomain is prescribed using the convergence rate and the slab dip; velocity in the lid is forced to equal zero. Within the wedge we solve the incompressible Stokes equation with no buoyancy term. There are no-slip conditions on the wedge–lid and wedge–slab boundaries, and there is a no-stress condition applied at the wedge inflow–outflow boundary. Temperature is fixed at zero on the top of the domain; the slab is given a temperature profile using the one-dimensional conducting slab solution; this solution is also applied to fix the temperatures on the inflow boundary. Complete details are provided in ref. 10.

The domain is 600 km by 600 km, oriented as shown in Fig. 1, with a mesh spacing of 1 km. We have checked convergence of our solutions by comparing them, for combinations of parameters that span the range of interest, against solutions with mesh spacing of 0.5 km. We find differences between the two meshes at the level of less than  $10^\circ \text{C}$  between temperatures calculated on the top of the slab or near the wedge corner.

An important parameter in thermal models of subduction zones is the depth at which the slab couples to the mantle wedge with a no-slip condition<sup>37</sup>. The maximum extent of shallowly dipping thrust faulting on the plate interface in subduction zones is

about 45 km (ref. 38) and, allowing for a transitional zone of aseismic sliding below this depth, we fix the depth to the top of full coupling between wedge and slab at 56 km.

**Measurement of depth to top of slab beneath arcs.** We apply the methods of ref. 8 to the 45 arc segments studied by ref. 13. For each arc segment we formed cross-sections of the seismicity and estimated  $D$  as described by ref. 8. We used the tables given in the supplementary information of ref. 13 to check whether the segment boundaries of the two studies agreed; except where noted in our Supplementary Table, the disagreements were minor and we used the segment boundaries of ref. 8.

We could not obtain reliable estimates of  $D$  for seven arcs, either because of the sparsity of earthquakes or because the volcanoes are spread out so widely that a clear volcanic front cannot be identified (see remarks on individual arcs, in the Supplementary Information). We combined the remaining volcanoes into 35 arc segments for which  $D$  estimated by the two different methods can be reconciled. An example of illustrating the two approaches, and their reconciliation in one of the more problematic arcs, can be seen in Supplementary Fig. 2. The Marianas, Lesser Antilles, and Vanuatu arcs are treated as single arcs here; each is split into northern and southern segments by ref. 41.

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## A DEPTH TO THE TOP OF THE SLAB BENEATH VOLCANIC ARCS

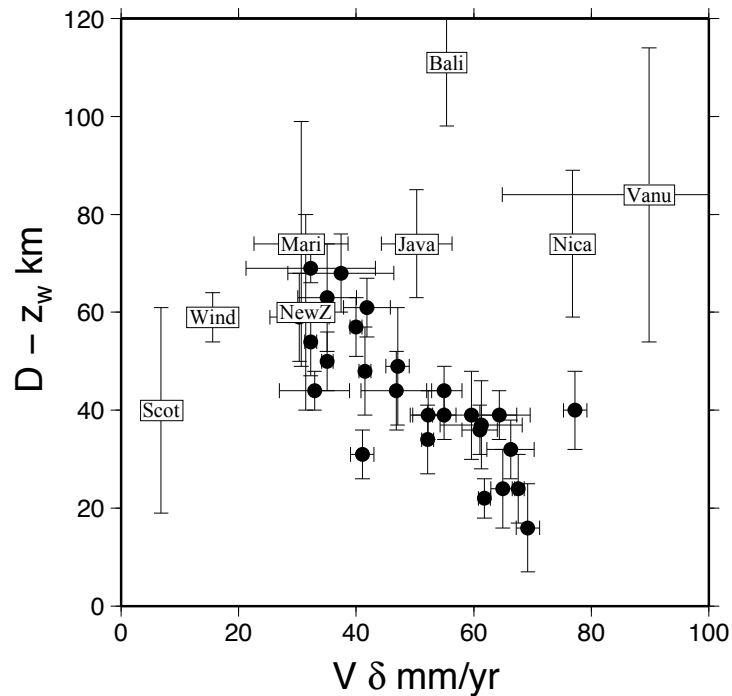
*England et al.* (2004) found a strong negative correlation between  $D$  and the descent speed of the slab for 22 volcanic arcs for which they could identify well-defined locations for the slab. That correlation was questioned by *Syracuse and Abers* (2006), who included in their study a half-dozen arc segments with steep ( $> 65^\circ$ ) dips, which *England et al.* had excluded. Here we revisit this issue, finding 35 arc segments for which  $D$  estimated by the two different sets of authors can be reconciled, and amplifying the reasons given by *England et al.* (2004) for excluding arcs for which they could not obtain reliable estimates of  $D$ .

The purpose of this exercise is to derive a set of estimates of  $D$  that may be used to test hypotheses about the relations between the thermal structure of arcs and the melting processes that take place beneath them. Several of the arcs that *Syracuse and Abers* (2006) identified as discordant with the trend identified by *England et al.* (2004) are shorter than 2-400km; for such segments, one should expect perturbations to the thermal regime, caused by flow around the edges of the slab (*Kincaid and Griffiths* 2003). In the case of the Philippine arcs, cold material from adjacent trenches may cool the wedge. For a few arcs, the distribution of seismicity makes determination of  $D$  problematic.

The total length of the arcs investigated by *England et al.* (2004) and *Syracuse and Abers* (2006) is about 30,000 km; there is approximately a further 6,000 km of arc below which the seismicity is absent, or too sparse for measurement of  $D$  by the methods employed by either set of authors. In the notes accompanying the Table (below) we identify the seven arcs (total length  $\sim 4000$  km) for which  $D$  was measured by *Syracuse and Abers* (2006), but which we have not included in our analysis, for one or more of the reasons given in the previous paragraph.

### A NOTE ON BACK-ARC SPREADING

As did *England et al.* (2004); *England and Wilkins* (2004), we use for the convergence rate  $V$  the rate of motion between the bounding plates, and do not include the rate of any back-arc spreading that occurs. Our rationale is based on the analysis that leads to the scaling relations developed in this paper. Temperatures in the mantle wedge are determined by the advection of heat by flow in the wedge. In subduction zones without back-arc spreading, all the flow is directed around the wedge corner, thus carrying heat beneath the arcs. When back-arc spreading occurs, there is greater net flow in the wedge, but the increased flow goes out through the back-arc spreading centre, rather than around the corner of the wedge (see *Ribe* 1989; *Conder et al.* 2002). Because the system is dominated by the advection of heat by the flow, it is reasonable to expect that any extra heat brought into the wedge by the back-arc spreading goes – with the additional flow – out through the back-arc, rather than contributing extra heat to the region below the arc. Therefore, the amount of heat advected *towards the arc* in a subduction zone with back-arc spreading should be comparable to the amount that would be carried in the absence of spreading. For this reason, we calculate  $V$  excluding the rate of back-arc spreading.



**Figure S1.** Figure corresponds to Figure 3a of the main text, except that arcs specifically discussed here are identified by their four-letter abbreviations (Table S1).  $z_w$ , the depth to the top of full coupling between wedge and slab, is taken to be 56 km (see Methods).

#### *Correlation between $D$ and $V\delta$*

In this paper we use the correlation between the depth to the top of the slab  $D$ , inferred from the distributions of intermediate-depth seismicity in the subduction zones and  $V\delta$ , the product of the convergence rate of the plates with the dip of the slab. We note in the body of the text that the rank-order correlation coefficient between  $D$  and  $V\delta$  is  $R = -0.34$  which, for a sample set consisting of 35 observations, has a probability of 2.6% of arising by chance. Because the distributions of  $D$  and  $V\delta$  are apparently non-normal as shown in Figure S2, we first establish that this probability is robust. We generated  $10^6$  random sampling sets of the data of Table S1, holding  $V\delta$  for each arc at its observed value, but shuffling the values of  $D$  and adding noise from a uniform distribution of width equal to the uncertainty in  $D$  in Table S1. The distribution of  $R$  for these samples is normal and, indeed, only 2.6% of these values are lower than  $R = -0.34$  (Figure S2c).

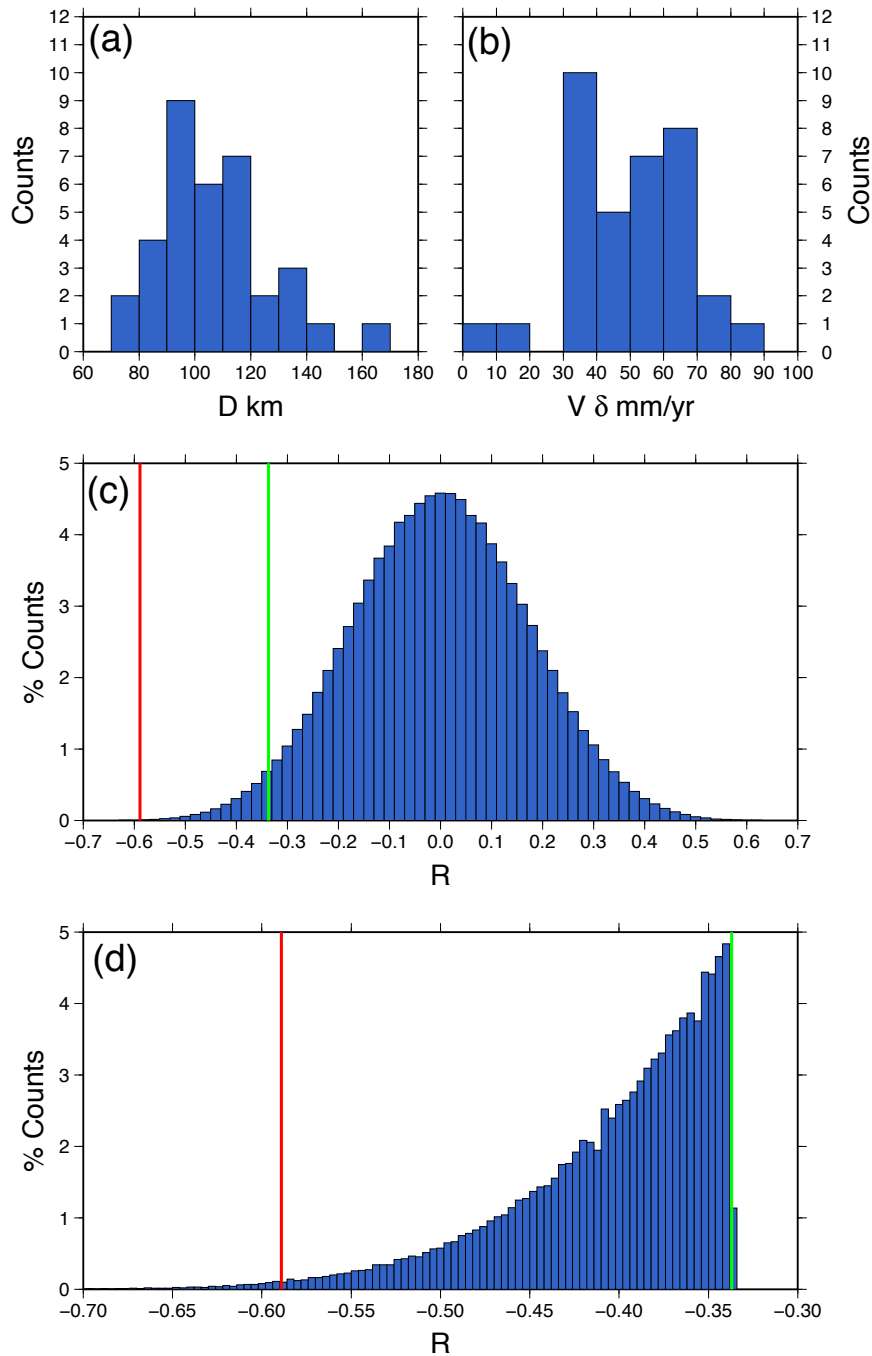
The significance of this correlation is sensitive to the presence of a small number of arcs that have large uncertainties in the measurement of  $D$ . To quantify this sensitivity, we exclude from the data set the 5 arcs for which the uncertainty in  $D$  is 15 km or greater. These are: New Zealand and Mariana, which lie on the general trend, and Vanuatu, Nicaragua, and Scotia, which lie off the trend (Figure S1). The rank-order coefficient for the remaining 30 arcs is -0.59

which, for a sample set containing 30 observations, has a probability of  $6 \times 10^{-4}$  of arising by chance. However, these are 30 observations chosen from among 35, with the criterion for exclusion being the magnitude of the uncertainty in  $D$ . It is reasonable to ask whether, by excluding these 5 arcs from the data set, we have artificially increased the significance of the correlation.

To address this question, we calculate the distribution of  $R$  obtained by randomly shuffling the data of Table S1 and removing from each sample the 5 observations with the largest rank-order difference (i.e. those observations that detract most from the significance of the correlation). We then recalculate  $R$  for each reduced sample of 30 arcs, and discard all samples with values of  $R$  is greater than  $-0.34$ ; the result is shown in Figure S2d. This procedure is designed to be more biased towards increasing the significance of  $R$  than our criterion of exclusion of arcs on the grounds of uncertainty in  $D$ . It yields a population of samples with correlation coefficients  $R$  that are at least as significant as that shown by the full data set. Only 2.3% of the samples have a value of  $R$  that is more negative than  $-0.59$ . This indicates that our exclusion of the arcs with high uncertainty in  $D$  has a low probability of artificially increasing the significance of the correlation.

Six segments of arc – Scotia, Java, Bali, Nicaragua, West Indies and Vanuatu, total length about 4,000 km – exhibit values of  $D$  that do not show the same correlation with descent speed seen for the other arcs. The analysis confirms the conclusion of *Syracuse and Abers* (2006) that some arcs – particularly those with slab dips above  $60^\circ$  – do not fit the trend observed by *England et al.* (2004). The analysis also reaffirms, however, the conclusion of *England et al.* (2004) that there is a significant correlation between  $D$  and the descent speed ( $\sim V\delta$ ) of the slab for the remaining 30 arcs, whose total length is about 22,000 km.





**Figure S2.** **Panel (a)** Distribution of measured values of depth to the top of the slab,  $D$ , beneath volcanic arcs (Table S1). **Panel (b)** Distribution of convergence rate  $V$  multiplied by slab dip  $\delta$  for these arcs. **Panel (c)** The distribution of rank-order coefficient  $R$  obtained by  $10^6$  random shufflings of the data of panels (a) and (b). The green line corresponds to  $R = -0.34$  and the red line to  $R = -0.59$  (see text). **Panel (d)** The distribution of rank-order coefficient  $R$  obtained by  $3 \times 10^7$  random shufflings of the data of panels (a) and (b), with the removal of 5 points from each set, as describe in the text. Only those realizations with  $R < -0.34$  are retained.

**Table 1.** The segments of arc considered in this study; they include all the segments considered by *England et al. (2004)* and by *Syracuse and Abers (2006)*; the two Vanuatu segments of Syracuse and Abers have been combined into one, as have the two Antilles (West Indies) and the Marianas. The first two columns give the segment names and values of  $D$  determined by *England et al. (2004)*; the third and the fourth give the segment names and values of  $D$  determined by *Syracuse and Abers (2006)*. The remaining columns give the names, estimates of  $D$ , and subduction parameters for all the segments studied here. Note that for several arcs, the disagreement in estimates of  $D$  between *England et al. (2004)* and *Syracuse and Abers (2006)* stems from differences in naming arc segments and/or different choices of segment boundaries. When only a portion of an arc segment of *Syracuse and Abers (2006)* has been used, the value of  $D$  quoted in our table is the average of their measurements of  $H$  (equivalent to our  $D$ ) for the different parts of the segment, not the value listed in their Table S1.

England et al., 2004		Syracuse and Abers, 2006		This Study						
Segment name	$D$ (km)	Segment name	$D$ (km)	Segment name	$D$ (km)	Dip, $\delta$ °	$V$ (mm/yr)	$V \sin \delta$ (mm/yr)	Note	
E. Aleutians	65	~C. Aleutians	72	E. Aleutians	72±9	61	65	57±2	1 Alee	
-	-	S. Kurile	78	S. Kurile	78±4	46	77	56±1	2 Kurs	
C. Aleutians	80	~W. Aleutians	78	C. Aleutians	80±8	61	61	53±2	3 Alec	
N. Kurile	80	N. Kurile	78	N. Kurile	80±7	49	79	60±1	4 Kurn	
E. Alaska Pen.	87	Alaska	85	Alaska	87±5	42	56	37±2	5 Alae	
Tonga	110	Tonga	88	Tonga	88±6	52	73	58±4	6 Tong	
-	-	Sunda Strait	90	Sunda Strait	90±7	49	61	46±1	7 Sund	
Kamchatka	95	Kamchatka	90	Kamchatka	92±5	46	76	55±3	8 Kamk	
Kermadec	95	Kermadec	91	Kermadec	93±9	65	54	49±7	9 Kerm	
S. Sumatra	100	S. Sumatra	89	S. Sumatra	95±9	56	61	51±10	10 Sums	
Guatemala	100	Guatemala	87	Guatemala	95±5	55	67	55±3	11 Guat	
W. Alaska Pen.	95	~E. Aleutians	85	W. Alaska Pen.	95±5	50	63	49±2	12 Alaw	
N. Japan	100	Central Honshu	95	Central Honshu	95±5	34	83	46±3	13 Chon	
-	-	Costa Rica	96	Costa Rica	96±8	59	75	64±2	14 Cost	
-	-	Scotia	96	Scotia	96±21	65	6	5±1	15 Scot	
W. Aleutians	100	~W. Aleutians	92	W. Aleutians	100±8	61	44	39±6	16 Alew	
S. Kuriles	100	Hokkaido	100	Hokkaido	100±5	42	75	52±3	17 Hokk	
-	-	Peru	100	Peru	100±4	29	65	31±4	18 Peru	
-	-	N. Honshu	104	N. Honshu	104±9	29	82	39±1	19 NHon	
Ryukyu	105	Kyushu	105	Kyushu	105±12	55	49	40±2	20 Kyus	
C. Chile	110	C. Chile	103	C. Chile	106±6	30	67	33±1	21 CChi	
-	-	N. Central Chile	110	N. Central Chile	110±7	24	77	31±1	22 NCCh	
-	-	N. Chile	113	N. Chile	113±6	29	79	39±1	23 NChi	
West Indies	115	N. & S. Antilles	108-122	West Indies	115±5	47	19	14±1	24 Wind	
N. Sumatra	115	C. Sumatra	101	C. Sumatra	115±9	37	47	28±5	25 CSum	
New Zealand	110	New Zealand	123	New Zealand	116±20	50	36	27±2	26 NewZ	
Izu	130	Izu	117	Izu	117±6	47	51	37±4	27 Izuu	
-	-	N. Sumatra	119	N. Sumatra	119±11	49	41	28±5	28 Sumn	
Bonin	120	Bonin	128	Bonin	124±8	63	24	28±4	29 Boni	
Aegean	125	-	-	Aegean	125±3	43	43	30±11	30 Aege	
-	-	Java	152	Java	130±11	40	72	46±6	31 Java	
-	-	Nicaragua	133	Nicaragua	130±15	62	71	63±1	32 Nica	
Marianas	115	N & S. Marianas	134-146	Marianas	130±25	65	27	25±8	33 Mari	
-	-	N. & S. Vanuatu	140	Vanuatu	140±30	78	66	69±25	34 Vanu	
-	-	Bali/Lombok	167	Bali/Lombok	167±13	46	69	50±2	35 Bali	
-	-	E. Banda Sea							36	
-	-	W. Banda Sea							37	
-	-	New Britain							38	
-	-	Solomons							39	
-	-	N. Philippines							40	
-	-	S. Philippines							41	
-	-	Central S. Chile							42	

*Notes to Table*

The designation “From *Syracuse and Abers* (2006)” indicates that we have used the procedure of *England et al.* (2004) to a segment defined by *Syracuse and Abers* (2006) (otherwise, we use the segment boundaries defined by *England et al.* (2004)). There is one exception to this procedure: in the case of N. Honshu, there are no entries in *Syracuse and Abers*’ supplementary information, so we accept their estimate of  $D$  (H in their notation). Uncertainties in  $D$  are given as in the original studies or, for those segments (the majority) that were studied by each set of authors, as the larger of the two uncertainties. Estimates of descent speed are from the respective original papers, except for Kyushu and the Marianas, as noted. This study was made before the publication of the MORVEL plate velocity model (*DeMets et al.* 2010), but we have checked that the figures we present would not be significantly affected if the newer velocities were to be employed.

(1) E. Aleutians: Central Aleutians of *Syracuse and Abers* (2006). *England et al.* (2004), Figure 3.

(2) S. Kurile: From *Syracuse and Abers* (2006). Contains southern part of North Kuriles and northern part of South Kuriles segments of *England et al.* (2004). Estimate of  $D$  by *England et al.* (2004), Figure 5, agrees with that of *Syracuse and Abers* (2006).

(3) C. Aleutians: Eastern part of West Aleutians of *Syracuse and Abers* (2006), between 181°E and 186°E. We use  $D$  of 80 km, picked from the cross-section of seismicity (*England et al.* (2004), Figure 4).

(4) N. Kurile: Estimate of  $D$  by *England et al.* (2004), Figure 5, agrees with that of *Syracuse and Abers* (2006).

(5) Alaska: Corresponds approximately to E. Alaskan Peninsula of *England et al.* (2004), Figure 6. We use the  $D$  from that study, but correct the dip, which was misprinted by *England et al.* (2004).

(6) Tonga: *England et al.* (2004), Figure 17, considered a shorter section of the Tonga arc than did *Syracuse and Abers* (2006). We have applied the procedure of *England et al.* (2004) to the longer segment of arc, and obtain the same value of  $D$  as do *Syracuse and Abers* (2006).

(7) Sunda Strait: From *Syracuse and Abers* (2006).

(8) Kamchatka: Estimate of  $D$  by *England et al.* (2004), Figure 8, agrees with that of *Syracuse and Abers* (2006).

(9) Kermadec: Estimate of  $D$  by *England et al.* (2004), Figure 9, differs by 4km from that of *Syracuse and Abers* (2006); we use the average of the two estimate of  $D$ .

(10) S. Sumatra: We use  $D$  of 100 km, picked from the cross-section of seismicity (*England et al.* (2004), Figure 13).

(11) Guatemala: Corresponds only to the Guatemala portion of the Guatemala-El Salvador segment of *England et al.* (2004), Figure 10. Estimate of  $D$  by *England et al.* (2004) agrees with estimates of *Syracuse and Abers* (2006), for the volcanoes of this section. The El Salvador

section is included with the Nicaragua section (below) because it appears to form part of the same, concave-towards-trench arc segment.

(12) W. Alaska Peninsula: East Aleutians of *Syracuse and Abers* (2006). We use  $D$  of 95 km, picked from the cross-section of seismicity (*England et al.* (2004), Figure 7). Average for  $D$  given by *Syracuse and Abers* (2006) over the same range of longitude is 90 km.

(13) Central Honshu: The Central and Northern Honshu arc segments of *Syracuse and Abers* (2006) correspond the N. Japan segment of *England et al.* (2004), Figure 11. We use the estimate of  $D$  by *England et al.* (2004).

(14) Costa Rica: From *Syracuse and Abers* (2006); *Syracuse et al.* (2008).

(15) Scotia: From *Syracuse and Abers* (2006). Note that this arc is short (about 400km), the slab dips steeply, and there is rapid back-arc spreading, so 3D effects (*Kincaid and Griffiths* 2003) may be important.

(16) W. Aleutians: Western part of West Aleutians of *Syracuse and Abers* (2006), between 176°E and 180°E. Shallower values of  $D$ , reported by *Syracuse and Abers* (2006) may reflect extrapolation of the location of top of slab based on seismicity further to the east. We use  $D$  of 100 km, picked from the cross-section of seismicity (*England et al.* (2004), Figure 4).

(17) Hokkaido: Corresponds to S. Kuriles segment of *England et al.* (2004), Figure 12. Estimate of  $D$  by *England et al.* (2004), agrees with that of *Syracuse and Abers* (2006).

(18) Peru: From *Syracuse and Abers* (2006).

(19) N. Honshu: From *Syracuse and Abers* (2006).

(20) Kyushu: Corresponds to Ryukyu segment of *England et al.* (2004), Figure 16. Estimate of  $D$  by *England et al.* (2004), agrees with that of *Syracuse and Abers* (2006). *Syracuse and Abers* (2006) use Philippine-South China angular velocity of *Sella et al.* (2002), which yields a descent speed of  $58 \pm 1$  mm/yr; we use, as do *England et al.* (2004) the Philippine-Eurasia angular velocity of *Seno et al.* (1993).

(21) Central Chile: The estimate of *England et al.* (2004), Figure 18, differs by 7 km from that of *Syracuse and Abers* (2006). We take the average of the two values of  $D$ .

(22) N. Central Chile: From *Syracuse and Abers* (2006).

(23) N. Chile: From *Syracuse and Abers* (2006).

(24) West Indies: We use  $D$  of 115 km, picked from the cross-section of seismicity (*England et al.* (2004), Figures 23 & 24). We combine the two segments.

(25) C. Sumatra: Seismicity defines a clear top to the slab at 115km beneath the volcanoes (*England et al.* (2004), Figure 22); depth given by *Syracuse and Abers* (2006) appears to be constrained by a small number of earthquakes at much shallower depth.

(26) New Zealand: The estimate of *England et al.* (2004), Figure 19, differs by 13 km from that of *Syracuse and Abers* (2006). We take the average of the two values of  $D$ .

(27) Izu: Re-examination of the cross-section of *England et al.* (2004), Figure 27, in the light of the analysis of *Syracuse and Abers* (2006) leads us to accept their estimate of  $D$ .

(28) N. Sumatra: From *Syracuse and Abers* (2006).

(29) Bonin: The estimate of *England et al.* (2004), Figure 25, differs by 8 km from that

of *Syracuse and Abers* (2006). We take the average of the two values of  $D$ . Dip, convergence rate, and descent speed taken from *Syracuse and Abers* (2006).

(30) Aegean: *England et al.* (2004), Figure 26.

(31) Java: The segment used by *Syracuse and Abers* (2006) is longer than that of *England et al.* (2004). We have applied the procedure of *England et al.* (2004) to the longer segment of arc, and obtain a value for  $D$  of 130 beneath the frontal volcanoes. This arc is much wider than most of the others considered here, and the value of 150 km for  $D$  given by *Syracuse and Abers* (2006), reflects an average for all volcanoes in the arc; their estimates of  $D$  for the volcanoes closest to the volcanic front lie in the range 130 to 140 km.

(32) Nicaragua: From Nicaragua and El Salvador of *Syracuse and Abers* (2006), who quote an uncertainty for  $D$  of 4 km for this segment. We base our value of 15 km for the uncertainty on the range in the depth to the top of the intermediate-depth seismicity in the 10-km band either side of the volcanic front, where the local dip of the top of the slab appears to be greater than  $70^\circ$ ; this range also appears to be consistent with the range of  $D$  given in the supplementary data of *Syracuse and Abers* (2006).

(33) Marianas: combines N. and S. Marianas of *Syracuse and Abers* (2006). Method of *England et al.* (2004) (cross-section shown in their Figure 20, and in Figure S3 here) indicates a minimum value for  $D$  of 115 km. Except in the southern-most Marianas, where relatively fewer earthquakes define the location of the slab beneath the arc, the method of *Syracuse and Abers* (2006) indicates that the volcanoes lie between the 100-km and 150-km contours of depth-to-slab (Figure S3). We assign  $D = 130 \pm 25$  km, which covers both estimates. *Syracuse and Abers* (2006) use Pacific-Philippine angular velocity of *Sella et al.* (2002), which yields a descent speed of  $17 \pm 1$  mm/yr; we use, as do *England et al.* (2004) the Pacific-Philippine angular velocity of *Seno et al.* (1993).

(34) Vanuatu: *Syracuse and Abers* (2006). Slab essentially vertical beneath arc.  $D$  could be anywhere between 110 km and 170 km; N and S. Vanuatu combined as single segment.

(35) Bali/Lombok: From *Syracuse and Abers* (2006).

(36) E. Banda Sea: *Syracuse and Abers* (2006). We do not include this segment, because it consists of about 250 km along-strike length of slab that is highly curved (radius of curvature about 200 km) giving a likelihood of 3D effects (*Kincaid and Griffiths* 2003); in addition, both  $D$  and the descent speed are poorly constrained.

(37) W. Banda Sea: From *Syracuse and Abers* (2006). Reliable intermediate-depth seismicity only E of  $122^\circ\text{E}$ , where we find  $D=115\text{km}$ . However, this section of arc is only 150km long, and we do not include it because of the likelihood of 3D effects (*Kincaid and Griffiths* 2003).

(38) New Britain: *Syracuse and Abers* (2006). Short arc segment, small radius of curvature; using techniques of *England et al.* (2004)  $D$  is uncertain.  $V$  uncertain because relative motions of bounding plates are not well constrained (*DeMets et al.* 2010). Not used in this study, although values for  $D$  and  $V\delta$  (*Syracuse and Abers* 2006) are consistent with the correlation described in this paper.

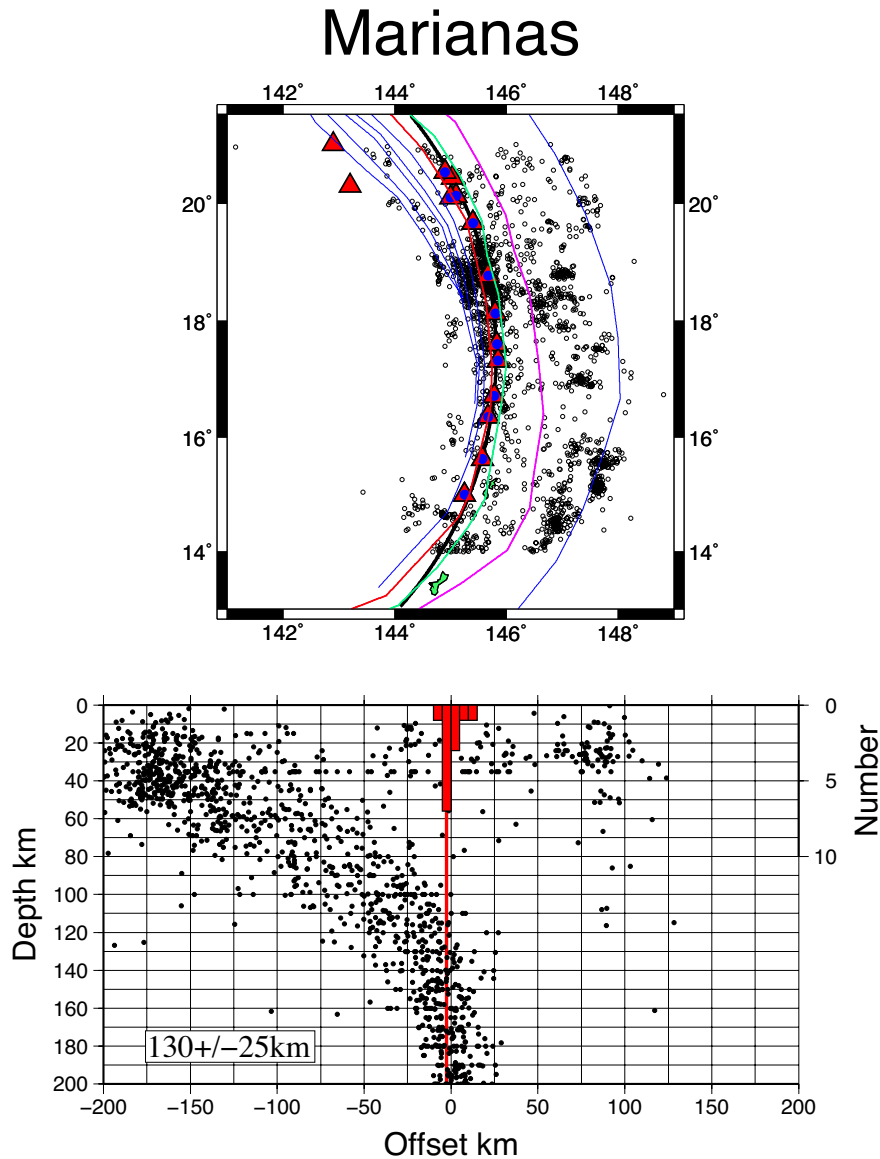
(39) Solomons: *Syracuse and Abers* (2006). Not used here; too few hypocenters beneath arc

for cross-section to be interpretable. Slab changes dip rapidly in region immediately beneath the arc; top of slab could be as shallow as 90km, or as deep as 150 km. Apart from one volcano, 200km to the north, and near a corner in the slab, the volcanoes define an arc less than 200km long. In addition,  $V$  uncertain because relative motions of bounding plates are not well constrained (*DeMets et al.* 2010).

(40) N. Philippines: *Syracuse and Abers* (2006). Arc too diffuse to obtain estimate of  $D$  using techniques of *England et al.* (2004). Different, reasonable, assumptions about which volcanoes define the arc front lead to small circles with  $D$  between 100 and 150 km. This result is also borne out by the depth-to-slab contours of *Syracuse and Abers* (2006). Additionally, convergence along the East Luzon Trough, less than 200 km to the east of the arc, presumably delivers cold material to the wedge, and produces a distinctive thermal structure that could not be analyzed with any simple model for subduction zones.

(41) S. Philippines: *Syracuse and Abers* (2006). A range of small circles defining the arc front can be found, depending upon the choice of volcanoes; the intermediate-depth seismicity dips steeply beneath these lines, and estimates of  $D$  lie between about 100 km and 180 km. As with the Northern Philippines, there may be delivery of cold material to the wedge by subduction of opposite vergence within 200–300 km of the arc.

(42) Central S. Chile: Segment of *Syracuse and Abers* (2006); we do not use this segment, because the very small number of earthquakes beneath these volcanoes makes it impossible to define accurately the location of the slab. Cross-section drawn using the procedure of *England et al.* (2004), suggests that the top of the slab may be as much as 20km deeper than the 80km *Syracuse and Abers* (2006), obtained by fitting a plane through the seismicity.



**Figure S3.** Map and composite cross-section of the distribution of earthquakes and volcanoes in the Mariana arc, following the methods described by *England et al.* (2004). In the upper panel, epicentres of earthquakes from the catalogue of *Engdahl et al.* (1998) are shown by open circles, and the locations of volcanoes from *Siebert and Simkin* (2002) are shown by triangles. The black line shows the best-fitting small circle through the locations of the volcanoes of the volcanic front, which are denoted by blue dots within their triangles. Coloured lines show contours of the depths to the top of the slab determined by *Syracuse and Abers* (2006). The easternmost contour corresponds to a depth of 6km; purple, green and red correspond, respectively, to 50, 100, and 150 km. Deeper contours, at intervals of 50km, are shown in blue. In the lower panel, the hypocentres of the earthquakes shown in the upper panel are plotted with the horizontal coordinate being the offset from the small circle. Histogram on the top of this panel (scale to right) indicates the numbers of volcanoes located within 5-km bins of distance from the small circle. Vertical line indicates chosen location of the volcanic front.

## B THERMAL CALCULATIONS

### B.1 Scaling

Figure 2 of the main text shows that the temperatures in our variable-viscosity simulations obey the same scaling relations that were derived for the isoviscous case. The preservation of the scaling from the isoviscous case considered by *England and Wilkins* (2004) is not unexpected. On dimensional grounds, the temperature can depend only upon  $V$ ,  $r$  and the thermal diffusivity  $\kappa$  (in the combination  $Vr/\kappa$ ), and upon dimensionless parameters. The only dimensionless parameter in the constant-viscosity problem is the dip  $\delta$  of the slab; the temperature-dependent problem has an additional dimensionless parameter representing the temperature-sensitivity of the rheology, but this parameter is constant for all calculations, so it will not enter into the scaling. Thus, for either problem, temperature should scale with  $Vr/\kappa$  and some power of the dip.

The scaling can be understood from the balance between advection and diffusion of heat near the wedge corner (*England and Wilkins* 2004). The scale for the advection of heat is set by the volume flux of mantle toward the wedge corner, which is proportional to  $Vr\delta$ , so we expect the maximum temperature in the wedge to increase with  $Vr\delta$ . But heat is also conducted out of the wedge corner, principally down the steep thermal gradient in the boundary layer on top of the slab. The thickness of this thermal boundary increases, for fixed  $V\delta$ , approximately in proportion to  $\delta$  (*England and Wilkins* 2004), the second power of  $\delta$  in the scaling therefore arises because for fixed  $Vr\delta$ , less heat is conducted out of the wedge corner at large dip than at small dip.

### B.2 Uncertainties in temperatures

In formalizing this model of the thermal structure of subduction zones we have made several simplifications and we do not expect that our calculations will precisely represent the temperature fields in real subduction zones. We have treated the dip of the slab as constant, at the value representative of the depth range between the base of the plate and 200 km, whereas in reality the dip shallows to  $10^\circ$ – $20^\circ$  in the upper parts of the plate boundary. As *England and Wilkins* (2004) and *Conder* (2005) show, reasonable assumptions about the dissipation of heat on the plate boundary (the subduction fault between the trench and the wedge corner) can lead to temperatures at the top of the slab that are  $\sim 100 - 200^\circ\text{C}$  higher than we calculate here. The details of the thermal structure of the plate boundary have, however, only a minor influence on temperatures in the wedge interior because advection in the wedge is the dominant control on the temperatures there (*England and Wilkins* 2004). We carried out experiments in which the coupling depth was 44 or 68 km, and found that these changes, respectively, raised and lowered peak temperatures in the wedge beneath the arcs by about  $50^\circ\text{C}$ . We also allowed the parameters  $E$  and  $A$  in the rheology (see Methods) to change from the values used here; changing  $E$  by  $\pm 10\%$  changed temperatures in the interior of the wedge



by about  $\pm 50^\circ\text{C}$ , and decrease (increase) in  $A$  by a factor of 2.5 caused increase (decrease) in wedge temperatures by  $\sim 50^\circ\text{C}$ .

Each of these sets of experiments yields the same form of scaling of temperature with the subduction parameters that we describe above. We believe, therefore, that our model captures the correlations between subduction parameters and thermal structure.

### B.3 Thermal Erosion by Melt

Here we estimate the potential for thermal erosion of the overlying lithosphere by nominally anhydrous magma rising out of the core of the wedge. We consider a 2D column of rock of width  $w$  and height  $h$  extending from the top of the nose of the mantle isotherm at temperature  $T_C$  upward to the surface of the Earth. Within the column is a mixture of fluid magma and solid mantle in local volume proportion  $\phi$ . Each of these two phases has a distinct internal energy per unit mass  $e_i$  and velocity  $\mathbf{v}_i$ , where  $i$  can be either  $f$  for the magma or  $m$  for the matrix. We assume that they both have a density  $\rho$ , a thermal conductivity  $k$ , and a local temperature  $T$ . We will be concerned with the changing energy balance within this column as magma enters it from below. An integral representation of this balance is

$$\frac{d}{dt} \int_V E dV = - \int_{\partial V} (\overline{\rho e \mathbf{v}} - k \nabla T) \cdot \mathbf{n} dS, \quad (1)$$

where  $E$  is the bulk internal energy per unit volume and overbars represent phase averages (i.e.  $\bar{q} = \phi q_f + (1 - \phi) q_m$ ). This equation states that changes in internal energy of the column are due to advection and diffusion of internal energy across its boundary. We have neglected contributions from kinetic and potential energy.

The mantle matrix has internal energy  $e_m = c_P(T - T_0)$  while the magma has  $e_f = c_P(T - T_0) + L$ , where  $c_P$  is the specific heat at constant pressure and  $L$  is the latent heat per unit mass. The bulk volumetric internal energy is thus

$$E = \rho c_P(T - T_0) + \rho \phi L. \quad (2)$$

We assume that heat and mass can enter and leave the column only through the bottom and top, and that there is no flux through the sides. Then the temperature in the column depends only on height, going from  $T_C$  at the bottom to  $T_S$  at the surface. We approximate this as a two-part linear profile,

$$T(z) = \begin{cases} T_C & 0 \leq z \leq z_M, \\ T_C + (T_C - T_S) \frac{z - z_M}{h - z_M} & z_M < z \leq h. \end{cases} \quad (3)$$

Without thermal erosion due to upwelling magma,  $z_M = 0$  and energy diffuses out of the column at the surface at a rate  $k(T_C - T_S)/h$ , which is exactly balanced by the heat flux into the base of the column from the hotter mantle below, such that the temperature profile is in steady state. The effect of magmatic thermal erosion will be to push  $z_M$  upward into the column.

We consider magma entering the column from below. Because the phase densities are equal and the column has a fixed volume, the upward flux of magma into the column must be balanced by a downward flux of solid mantle out of the column. These occur at the same temperature, so the net advective heat transport into the column is  $ML$ , where  $M$  is the rate of flow of magma into the column (mass per unit length per unit time). Finally, assuming that the influx of magma into the column is balanced by freezing within the column, we can take the average porosity within the column to be constant.

Applying all of these assumptions to equation (1), we can rewrite it as

$$\frac{1}{2}w\rho c_P\Delta T\frac{dz_M}{dt} = ML - \frac{kw\Delta T}{h - z_M} + \frac{kw\Delta T}{h}, \quad (4)$$

where  $\Delta T = T_C - T_S$ . This equation states that the rate of thermal erosion within the column is controlled by the flux of magma into the base of the column, the diffusion of heat out through the upper surface, and the supply of heat from the mantle below. The last term in (4) ensures that  $z_M = 0$  is a solution when  $M = 0$ . Scaling  $z$  with  $h$  and  $t$  with the characteristic thermal diffusion time  $h^2\rho c_P/k$  gives the non-dimensional equation

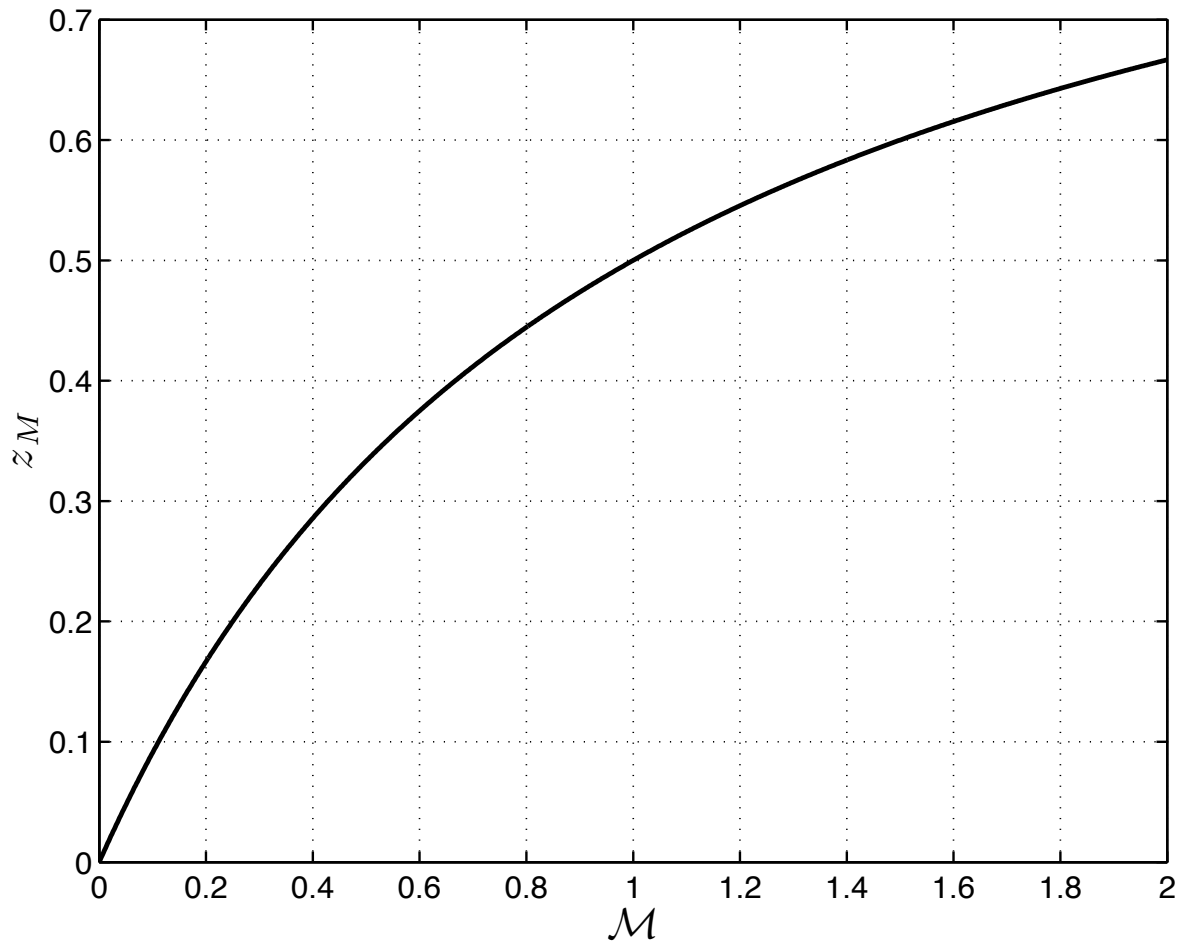
$$\dot{z}_M = 2 \left( \mathcal{M} - \frac{z_M}{1 - z_M} \right), \quad (5)$$

where  $\mathcal{M} = MLh/(kw\Delta T)$ , a dot above  $z_M$  indicates the time derivative, and  $z_M$  and  $t$  are now dimensionless variables. In steady state,  $z_M = \mathcal{M}/(\mathcal{M} + 1)$ ; this relationship is plotted in Figure S4.

We may make reasonably well constrained estimates of the quantities contained in  $\mathcal{M}$ . According to our analysis,  $H$  lies in the range 60–100 km (Figure 3b). We estimate the width of the column,  $w$ , from the width of the volcanic fronts, which is 10–15 km (*England et al.* 2004), so use  $h/w \sim 5$ . We take  $\Delta T = 1200$  K,  $k = 4$  W m<sup>-1</sup> K<sup>-1</sup> and  $L = 4 \times 10^5$  J kg<sup>-1</sup>.

We first estimate the likely magnitude of  $M$  for nominally anhydrous melting alone; calculation of dry melt production based on thermal models for subduction zones (*Conder et al.* 2002) yields an estimate of  $\sim 2 \times 10^{-3}$  kg m<sup>-1</sup> sec<sup>-1</sup> for the total rate of dry melt production in the wedge, which we may equate to  $M$  if we assume that all of this melt passes through the column. Using this estimate gives  $\mathcal{M} \sim 0.8$ , so this degree of dry melt production can thermally erode the overlying temperature structure so that the 1200°C isotherm is advected upward by over 40% of its original depth. Melting at 1/4 of this rate would raise the 1200°C isotherm by 20%, or about 10 km. Note that these estimates ignore the influence of along-arc focusing of the melt toward individual volcanic centres (e.g. *d'Ars et al.* 1995).

Estimates of the total melt production rate at arcs vary considerably, but are generally higher than the above estimate by more than a factor of two (e.g. *Crisp* 1984; *Reymer and Schubert* 1984; *Dimalanta et al.* 2002; *White et al.* 2006), so advection of heat by hydrous melt alone would be capable of raising the 1200°C isotherm by over 60% of its original depth. An equivalent effect would occur if less than half the heat advected by the hydrous melt were added to that carried by the dry melt.



**Figure S4.** A plot of the steady-state, dimensionless height of thermal erosion distance as a function of the dimensionless control parameter,  $\mathcal{M}$ . The curve is given by  $z_M = \mathcal{M}/(\mathcal{M} + 1)$ .

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